

Nuclear Physics from Lattice QCD

Kostas Orginos William & Mary -- JLab



Thomas Jefferson National Accelerator Facility

Collaborators



Un. of New Hampshire

Un. of Washington (Seattle)

College of William and Mary

Un. of Barcelona

LLNL

Motivation-Overview

- Study Multi-hadron systems
- Applications:
 - Nuclear physics (**A=2**, 3, 4 ...)
 - Spectroscopy (excited states, multi-quark hadrons etc.)
- Methodology: Use finite volume
 - Extract scattering phase shifts
 - Extract multi-hadron interaction properties
- Results and Conclusions



Charmonium Spectrum



Hadron Scattering

- Scattering processes from Lattice QCD are not straight forward
- Miani-Testa no-go theorem ('90) [and C. Michael '89]
- Infinite Volume:



Minkowski

- Finite volume: discrete spectrum
 - Avoids Miani-Testa no-go theorem [M. Luscher]

Scattering on the Lattice Luscher

Scattering amplitude:

At

finite volume one can show:
$$\Delta E_n \equiv E_n - 2m = 2\sqrt{p_n^2 + m^2} - 2m$$

$$P_{n} \text{ solutions of:}$$

$$p \cot \delta(p) = \frac{1}{\pi L} \mathbf{S} \left(\frac{p^{2} L^{2}}{4\pi^{2}} \right) \qquad \mathbf{S}(\eta) \equiv \sum_{\mathbf{j}}^{|\mathbf{j}| < \Lambda} \frac{1}{|\mathbf{j}|^{2} - \eta} - 4\pi\Lambda$$

Effective range expansion:

$$p \cot \delta(p) = \frac{1}{a} + \frac{1}{2}rp^2 + \dots$$

a is the scattering length

Luscher Formula

Energy level shift in finite volume:

$$\Delta E_n \equiv E_n - 2m = 2\sqrt{p_n^2 + m^2} - 2m$$

$$p_{n} \text{ solutions of:}$$

$$p \cot \delta(p) = \frac{1}{\pi L} \mathbf{S} \left(\frac{p^{2}L^{2}}{4\pi^{2}} \right) \qquad \mathbf{S}(\eta) \equiv \sum_{j=1}^{|\mathbf{j}| < \Lambda} \frac{1}{|\mathbf{j}|^{2} - \eta} - 4\pi\Lambda$$

$$p_{n} \cot \delta(p_{n}) = \frac{1}{a} + \cdots \qquad \frac{1}{a} = \frac{1}{\pi L} S \left(\frac{p_{0}^{2}L^{2}}{4\pi^{2}} \right) + \cdots$$

Expansion at p~0 :

$$\Delta E_0 = -\frac{4\pi a}{mL^3} \left[1 + c_1 \frac{a}{L} + c_2 \left(\frac{a}{L}\right)^2 \right] + \mathcal{O}\left(\frac{1}{L^6}\right)$$

a is the scattering length

c₁ and **c**₂ are universal constants

Technicalities

- How well can we compute energy levels in LQCD?
 - Ground state:
 - Relatively easy for mesons
 - Baryons are hard and get harder as their number grows
 - Excited states are more demanding but methods do exist
- Theoretical background
 - Two body problems are well understood [Lucher]
 - Multi-mesons recently done
 - [Beane et.al. PRD76;074507, 2007, Detmold et.al. PRD77:057502,2008]
 - Coupled channels: Works begin to appear [Lage et.al. arXiv:0905.0069]
 - More than two baryons: Work in progress.... [Luu <u>arXiv:0810.2331v1</u>]
 - Still a lot needs to be done!



 $m_{\pi} a_2 = -0.04196(12)$

 $m_{\pi} a_2 = -0.04223(28)$

- Physical point -- Lighter three point fit:
- Physical point -- Lighter two point fit:
- Physical point -- Quadratic fit (higher order): $m_{\pi} a_2 = -0.0426(4)$



 $m_{\pi} a_2 = -0.04196(12)$

 $m_{\pi} a_2 = -0.04223(28)$

- Physical point -- Lighter three point fit:
- Physical point -- Lighter two point fit:
- Physical point -- Quadratic fit (higher order): $m_{\pi} a_2 = -0.0426(4)$



Physical point -- Lighter three point fit: m_π a₂ = -0.04196(12)
Physical point -- Lighter two point fit: m_π a₂ = -0.04223(28)
Physical point -- Quadratic fit (higher order): m_π a₂ = -0.0426(4)



 $m_{\pi} a_2 = -0.04196(12)$

 $m_{\pi} a_2 = -0.04223(28)$

- Physical point -- Lighter three point fit:
- Physical point -- Lighter two point fit:
- Physical point -- Quadratic fit (higher order): $m_{\pi} a_2 = -0.0426(4)$



 $m_{\pi} a_2 = -0.04196(12)$

 $m_{\pi} a_2 = -0.04223(28)$

- Physical point -- Lighter three point fit:
- Physical point -- Lighter two point fit:
- Physical point -- Quadratic fit (higher order): $m_{\pi} a_2 = -0.0426(4)$

World results on I=2 π-π scattering lengths



Three meson interaction

NPLQCD: Phys.Rev.D78 (054514,014507) 2008



- Three pion interaction is non-zero
- Three kaon interaction vanishes

Nucleon-Nucleon

NPLQCD: Phys.Rev.Lett.97 2006



BBSvK: Beane Bedaque Savage van Kolck '02 W: Weinberg '90;Weingberg '91; Ordonez et.al '95

Fukugita et al. '95: Quenched heavy pions

Signal to Noise ratio for correlation functions (mesons)

 $C(t) = \langle m(t)\bar{m}(0) \rangle \sim Ae^{-M_m t}$

 $var(C(t)) = \langle m\bar{m}(t)m\bar{m}(0) \rangle \sim Ae^{-2M_m t} + Be^{-2M_\pi t}$

$$StoN = \frac{C(t)}{\sqrt{var(C(t))}} \sim Ae^{-(M_m - M_\pi)t}$$

• For pseudo-scalar mesons the signal does not deteriorate at large time

Signal to Noise ratio for correlation functions (baryons)

 $C(t) = \langle N(t)\bar{N}(0)\rangle \sim Ee^{-M_N t}$

 $var(C(t)) = \langle N\bar{N}(t)N\bar{N}(0)\rangle \sim Ae^{-2M_Nt} + Be^{-3m_\pi t}$

$$StoN = \frac{C(t)}{\sqrt{var(C(t))}} = \sim Ae^{-(M_N - 3/2m_\pi)t}$$

- The signal to noise ratio drops exponentially with time
- The signal to noise ratio drops exponentially with decreasing pion mass
- For two baryons: $StoN(2N) = StoN(1N)^2$

High Statistics runs

- Improved Wilson fermions on anisotropic lattices
- Single pion mass: 390MeV
- Very high statistics (~300x10³ correlators)
- Goals
 - Get clean signals
 - Investigate methods of extracting masses from correlation functions
 - Check feasibility of 3 body calculations

Signal to Noise Ratio



- Exponential loss of signal at large time
- Slopes become larger as baryon number increases

Exponential Slope of the variance

 $var \sim C_N e^{-2AM_N t} + C_\pi e^{-3AM_\pi t}$



Energy Shifts



pion mass: 390MeV

Scattering Lengths



pion mass: 390MeV



- Errors on scattering nucleon-nucleon scattering length as function of computational resources
- Only cost for correlation function calculation presented
- Includes expected speedup from eigCG [Stathopoulos, KO 2007]

Summary

- Mesonic (pseudo-scalar) many body systems are successfully studied in LQCD
 - Scattering lengths
 - Three body interaction
- Baryons are hard
 - Statistical noise require high statistics calculations
- Advent of peta-scale computing together with new algorithmic developments will provide interesting results over the next few years